

# Concepts in Condensed Matter Physics: Exercise 4

Due date: 14/07/24

## 1 The XY – sine-Gordon duality and the BKT critical behavior

In this question you are asked to show the equivalence between the XY model and the sine-Gordon model:

$$S_{SG} = \frac{c}{2} \int d^2x (\nabla\theta)^2 - g \int d^2x \cos\theta \quad (1)$$

where  $\theta$  is a non-compact real scalar field, and re-derive the RG equations near the BKT transition.

1. Expand  $Z_{SG} = \int D\theta e^{-S_{SG}}$  in powers of  $g$  explicitly and show that it has the form

$$Z_{SG} = \sum_{n=0}^{\infty} \frac{(g/2)^{2n}}{(n!)^2} \prod_{j=1}^{2n} \int d^2x_j \left\langle \exp \left( i \sum_{j=1}^{2n} (-1)^j \theta(x_j) \right) \right\rangle, \quad (2)$$

Where the brackets denote averaging with the free part  $S_0 = \frac{c}{2} \int d^2x (\nabla\theta)^2$ . Hint: recall that the free part is translationally invariant such that  $\left\langle \left( \prod_{a=1}^N e^{i\theta(x_a)} \right) \left( \prod_{b=N+1}^{N+M} e^{-i\theta(x_b)} \right) \right\rangle$  is non-zero only for  $N = M$ .

2. Using the properties of the Gaussian average, namely  $\langle e^A \rangle = e^{\frac{1}{2}\langle A^2 \rangle}$  for  $A$  which is a linear combination of the field  $\theta$  and the following identity

$$\left\langle (\theta(x) - \theta(x'))^2 \right\rangle = \frac{C(x-x')}{c} = \frac{1}{2\pi c} \log \left| \frac{x-x'}{\xi} \right|, \quad (3)$$

show that the partition function may be written as follows

$$Z_{SG} = \sum_{n=0}^{\infty} \frac{(g/2)^{2n}}{(n!)^2} \prod_{j=1}^{2n} \int d^2x_j \exp \left( \frac{1}{2c} \sum_{j<i}^{2n} \sigma_i \sigma_j C(x_i - x_j) \right), \quad (4)$$

where  $\sigma_i$  denotes the sign of the vortex and  $\xi$  is a short length cutoff. This is exactly the partition function of the Coulomb gas obtained in class.

3. Repeat the derivation of the RG differential equations near the BKT transition, namely

$$\frac{dy}{d\ell} = xy, \quad \frac{dx}{d\ell} = y^2. \quad (5)$$

What are  $x$  and  $y$  in terms of  $c$  and  $g$ ? Here  $\ell = \log \frac{r}{\xi}$ .

4. Use the above equations to determine the screening length  $\xi_+$  on the disordered side close to the transition. Do this by estimating the value of the running parameter  $\ell$  at which  $x$  and  $y$  reach order 1. Explain physically why  $\xi_+$  is the screening length.
5. Obtain the superfluid stiffness  $J$  as a function of  $t = T - T_c$  and show that it has a universal jump at  $T_c$ .

## 2 Self consistent Harmonic approximation

In this question, we will do a variational approximation to estimate the critical temperature, and the scaling of the gap in the sine-Gordon (SG) model. We will show that an upper-bound of the free energy can be found by taking expectation values with respect to the gaussian action

$$S_{\text{var}} = \int d^2x \left[ \frac{c}{2} (\nabla\theta)^2 + m^2\theta^2 \right] \quad (6)$$

$$Z_0 = \int \mathcal{D}\theta e^{-S_{\text{var}}}, \quad (7)$$

where  $m$  is a variational parameter, and  $c$  is the same as in the SG model. We can write the exact expression

$$Z_{\text{SG}} = \int \mathcal{D}\theta e^{-S_{\text{SG}}} = \int \mathcal{D}\theta e^{-S_{\text{var}}} [e^{S_{\text{var}} - S_{\text{SG}}}] = Z_0 \langle e^{S_{\text{var}} - S_{\text{SG}}} \rangle_{\text{var}}, \quad (8)$$

and approximate it by

$$Z_{\text{SG}} \approx Z_{\text{var}} = Z_0 e^{\langle S_{\text{var}} - S_{\text{SG}} \rangle_{\text{var}}}. \quad (9)$$

Next, we define the exact and approximated Free energies

$$F_{\text{SG}} = -\frac{1}{\beta} \log(Z_{\text{SG}}) \quad (10)$$

$$F_{\text{var}} = -\frac{1}{\beta} \log(Z_{\text{var}}) = F_0 + \frac{1}{\beta} \langle S_{\text{SG}} - S_{\text{Var}} \rangle_{\text{var}} \quad (11)$$

1. Show that  $F_{\text{var}}$  is strictly an upper bound for  $F_{\text{SG}}$ . Can you relate this fact to the definition of the free energy and of the canonical ensemble?
2. Compute  $\partial(\beta F_{\text{var}})/\partial m^2$  and find  $m$  that minimizes  $F_{\text{var}}$ . (You may have to introduce an ultraviolet cutoff  $\Lambda$  to the calculation)
3. How can we identify the phase transition point from this calculation? hint: How does the correlator  $\langle e^{i\theta(r)} e^{-i\theta(0)} \rangle$  decay for finite  $m$ ? How does it decay for  $m \rightarrow 0$
4. Take  $\Lambda \rightarrow \infty$  and find the critical value of  $c$ . Show that  $m \propto \Lambda^\alpha$  next to the transition, and find  $\alpha$ . How does this compare to the results from RG analysis of the BKT transition?

### 3 RG analysis for a localized perturbation

In this question, we will analyze the effect of local terms on a 1D superfluid. Start from a one-dimensional quantum system described by the following Ginzburg-Landau action:

$$S = \int_0^{1/T} d\tau \int dx \psi^* \left( \partial_\tau - \frac{\nabla^2}{2m} + \alpha + \beta |\psi|^2 \right) \psi. \quad (12)$$

1. Write the complex field  $\psi$  in a polar representation,  $\psi = \sqrt{\rho} e^{i\theta}$ . Write the corresponding action in terms of the polar representation. What is the mean-field value of  $\rho$  (denoted by  $\rho_0$ )? Show that the action has a  $U(1)$  symmetry acting as  $\theta \rightarrow \theta + \text{const}$ .
2. Since the above is a continuous symmetry, the  $\theta$ -field is a Goldstone mode. Write  $\rho = \rho_0 + \delta\rho$ , expand the action to second order in  $\delta\rho$  and integrate out the massive Gaussian fluctuations  $\delta\rho$ . Neglecting higher orders, show that the action of this Goldstone mode then takes the form

$$S_{XY} = -\frac{1}{2} \int_0^\beta d\tau \int dx \theta (\chi \partial_\tau^2 + \rho_s \nabla^2) \theta. \quad (13)$$

This action describes the low energy excitations of a superfluid. Go to zero temperature, and rescale  $x, \tau$  by defining  $\tau = \tilde{\tau}/\sqrt{\chi}$ ,  $x = \tilde{x}/\sqrt{\rho_s}$  such that the action becomes rotationally symmetric in  $x - \tau$  plane.

3. Add a term of the form

$$S_1 = -u \int_0^\infty d\tilde{\tau} \cos \theta(\tilde{x} = 0, \tilde{\tau}) \quad (14)$$

to the action. Integrate out the fields at  $\tilde{x} \neq 0$ . Hint: you can do this for  $u = 0$  and add a finite  $u$  in the end. Write an effective action for the field  $\theta(\tilde{x} = 0, \tilde{\tau})$ . Show that the effective action is given by:  $S_{eff} = S_0 + S_1$  with  $S_0 = r \int d\omega |\omega| \theta(0, \omega) \theta(0, -\omega)$ . What is the value of  $r$ ? Hint: you can find that by calculating  $\langle \theta(0, \omega) \theta(0, -\omega) \rangle$  in the original model and the effective model. Use the relation  $\theta(x, \tau) = \left(\frac{1}{\sqrt{2\pi}}\right)^2 \int dk d\omega \theta(k, \omega) e^{ikx + i\omega\tau}$ .

4. Using the effective theory obtain the RG equation for  $u$  (to first order).
5. Draw the flow diagram.